the temperature predicted for the onset of superfluidity.8 The researchers found that the gap increased as the temperature fell. Furthermore, they noted that the energy gap opened well above the predicted temperature; they cited that behavior as evidence that the pairs were correlated at higher temperature but became coherent only below  $T_a$ .

Earlier this year. Thomas and his group at Duke teamed with Kathryn Levin and her coworkers from the University of Chicago to undertake a thermodynamic measurement of a fermion cloud.9 The collaborators determined the heat capacity of the Fermi gas in the strong-interaction region and reported an abrupt change in slope, which they attributed to the formation of a superfluid. The transition temperature is close to the predicted onset of superfluidity. The collaborators had to devise a modeldependent calibration of the temperature in the strong-interaction region, where the well-understood relation between the shape of the atomic cloud and its temperature breaks down.

More recently, the Duke experimenters have reported that the rate of damping of a radial oscillation in a trapped cloud of gas slows near the same temperature at which they saw the change in heat capacity.<sup>10</sup>

Grimm recognizes that his group's experiment,8 as well as that of the Duke-Chicago team, "needs some theoretical support for a conclusion on superfluidity," though, he adds, such support is available in both cases.

In a forthcoming paper, Hulet and his Rice University coworkers have studied further the nature of fermionic pairs.11 They have measured the pair correlation as the interaction strength increases through the Feshbach resonance and, for the first time, into the weakly interacting BCS regime. They find that the pair correlation varies smoothly across this region, growing smaller as the size of the pairs grows. The data show a deviation from two-body physics in the resonance region and imply that manybody physics must be invoked there.

#### Still more horizons

Theorist Jason Ho of Ohio State University is very excited about exploiting the atomic systems to learn more about unusual many-body phenomena similar to those in high- $T_c$  superconductors. In high-T<sub>c</sub> materials, he says, the normal state is not the usual Fermi liquid, but instead "is very weird." One might be able to study the normal state of the superfluid atomic system by putting in so much angular momentum that the superfluidity will go away and reveal the normal state. It will also be interesting to measure many other properties of the superfluid system, but doing so will require model-independent methods of thermometry.

Ketterle sees the vortices as a way to map out superfluidity in a broader range of interaction strengths. There are new questions to answer, and theoretical predictions are already emerging about such things as the structure of vortex lines and the density depletion in the cores.

#### Barbara Goss Levi

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# Single, Physics-Based Model Accounts for the Mechanical Properties of Diverse Biopolymer Gels

Despite their variety of composition, stiffness, and strength, the proteinbased filaments that support and connect cells behave in the same, predictable way when strained.

During each heartbeat, the aorta's walls expand to hold the pulse of blood, then rebound elastically to propel the blood downstream. At rest, a healthy human heart has a systolic blood pressure of about 16 kilopascals. But for strenuous, impulsive effort, like dashing upstairs or hoisting a sack of mulch, the systolic pressure can double, even triple.

If the aorta's walls responded linearly to such increases, they'd balloon and most likely burst. Fortunately, Nature has endowed the aorta with a nonlinear property called strain stiffening. As stress increases so too does the ratio of stress to strain, the elastic modulus.

The aorta's strain stiffening arises

from its structure, a complex set of concentric, alternating layers of rubbery and stiff tissues. It's surprising, therefore, to find strain stiffening in the far simpler gels that form the support structures within and between cells (see figure 1).

Those gels consist of weblike filamentary networks of biopolymer embedded in the ambient extracellular or intracellular solution. They perform different functions and bear different stresses and strains. Yet remarkably, despite the variety of composition, strength, and flexibility, the gels' stress-strain curves follow the same basic shape.

Physicists, when they encounter such similar behavior, look for a uni-

versal equation or model. And that's what Cornelis Storm of the University of Leiden and his collaborators have now found. According to their analysis, a biopolymer gel's nonlinear response to strain depends principally on the stiffness of its constituent filaments and the density of the links between them.1 The work suggests that human attempts to mimic certain biomaterials might turn out to be easier than expected.

Storm's paper builds largely on the earlier work of two of his coauthors, theorist Fred MacKintosh of the Free University of Amsterdam, and experimenter Paul Janmey of the University of Pennsylvania. The other coauthors are also from the University of Pennsylvania: Tom Lubensky, who helped out with the theory, and Jennifer Pastore, who made the measurements the model sought to explain.

## Semiflexible polymers

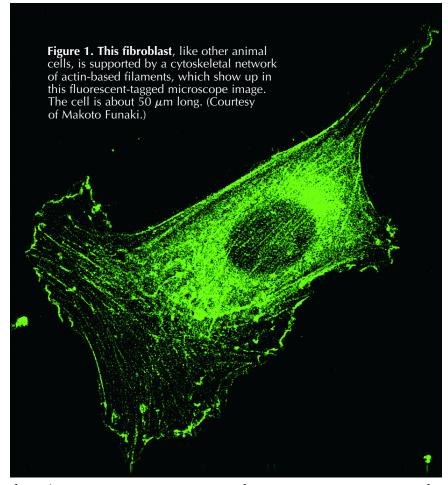
By the 1980s, biologists had identified the main network-forming proteins that animal cells rely on for structural support: collagen and fibrin outside cells and cytoskeletal polymers such as actin inside cells. Close up, as figure 2 shows, the networks resemble spaghetti spread out on a plate.

At first, the need for, or even aptness of, a general model for the gels' mechanical properties wasn't clear. Actin is much stiffer than collagen or fibrin and could, in principle, be treated as a mesh of semirigid girders. Fibrin, on the other hand, can be stretched to several times its original length before breaking.

Despite the range of stiffness, biopolymers are all what rheologists call semiflexible. That is, the length scales on which filaments undulate thermally and form links with each other are roughly the same. By contrast, the isoprene polymers that constitute rubber wiggle on a scale much shorter than the typical gap between crosslinks.

Semiflexibility was the starting point for a general, physics-based model that MacKintosh, Janmey, and Josef Käs developed ten years ago.<sup>2</sup> Their aim was to account for the large shear modulus and nonlinear response of biopolymer gels.

In the model, the origin of a filament's elasticity is entropic. Unstrained, a filament can adopt one of many wavy configurations. Stretching it, however, reduces the number of accessible configurations. Like a writhing snake pulled straight by its head and tail, the filament resists extension. The stiffer the filament, or the stronger the snake, the greater



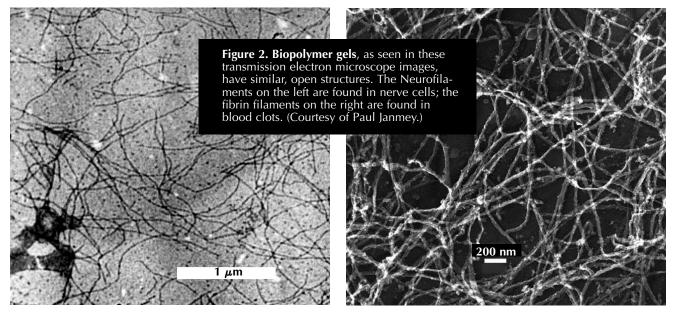
the resistance.

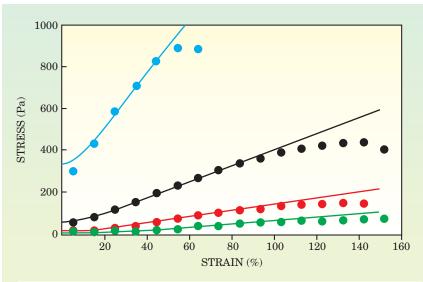
To predict the elastic modulus of a network of such filaments, MacKintosh, Janmey, and Käs assumed the network looked, on average, like an open, 3D mesh. That approach could, indeed, account for the low-strain elastic modulus, but, as is often the case, the need for a better model arose

when new measurements appeared.

## Going nonlinear

Reconstituting biopolymers in vitro is not easy. The protein monomers have to be isolated, purified, and then coaxed into polymerizing without the help of the living cells' complement of specialized enzymes. Measuring the





**Figure 3. The stress–strain behavior** of biopolymer gels depends on the concentration of the biopolymer filaments. Here, the data points represent measurements of reconstituted fibrin gels at four different concentrations: 0.5 mg/ml (green), 1.0 mg/ml (red), 2.0 mg/ml (black), and 4.5 mg/ml (blue). As the curves show, the model developed by Cornelis Storm and his coworkers provides a good fit to the data over a range of strains and concentrations. (Adapted from ref. 1.)

stress-strain curves of the gels involves placing the samples in a rheometer—essentially two parallel plates, one of which is movable. The rheometer records the stress required to maintain a fixed strain.

In 2001, Janmey and Pastore extended their measurements of seven different biopolymer gels into the nonlinear, strain-stiffening regime to examine the effects of varying filament stiffness and mesh size. At that time, Storm was a postdoc in Janmey's department. When he saw the data, he wondered how the earlier model could be extended.

He retained the entropic component of the filament stiffness and added a term to account for the elasticity of a fully straightened filament. To model the network, Storm assumed the filaments are distributed isotropically and randomly. Deriving the bulk modulus then boiled down to integrating across the network's nodes.

That step proved mathematically messy, until Lubensky discovered a fortuitous simplification. If the local deformation matches the bulk deformation—that is, if the deformations are affine, as had been assumed—then certain symmetries lead to a compact, practical expression for the stress tensor.

On the scale of a cell, as figure 1 shows, or in low-density regions where inhomogeneities become more important, the affine assumption

breaks down. Still, as figure 3 shows, the model proved to account for the data remarkably well. And although the final formula for the bulk modulus has to be solved numerically, it contains just one independent parameter.

It's tempting to attribute such apparently universal behavior to a basic requirement identified and picked out by evolution. But, cautions Janmey, that might not be the case for the strain stiffening of biopolymer gels. Some cells, such as long, delicate neurons, do need to withstand high strain, but cells also need to avoid obstructing the free flow of large molecules and even larger organelles. Strain stiffening of biopolymer gels could, perhaps, be an accidental byproduct of their open construction.

#### **Charles Day**

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